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( $\text{Ni}_{1-x}\text{Cu}_x$ ) $_{80}\text{P}_{20}$  METALLIC GLASSES

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## АННОТАЦИЯ

При комнатной температуре изучались ЯМР-параметры ядер  $^{31}\text{P}$  как функции соотношения Ni/Cu в аморфных сплавах  $(\text{Ni}_{1-x}\text{Cu}_x)_{80}\text{P}_{20}$  ( $0 \leq x \leq 0,77$ ), полученных быстрым охлаждением из расплава. Измеренные полнезависящие вторые моменты совпадают с полученными ранее результатами модельных расчетов для дипольного расширения. Найдено, что дипольное взаимодействие P-P не зависит от содержания Cu. Даются некоторые данные по сдвигу Найта и времени спин-решеточной релаксации и их краткое обсуждение.

## KIVONAT

$^{31}\text{P}$  NMR vizsgálatot végeztünk szobahőmérsékleten a Ni/Cu arány függvényében olvadákból gyors hűtéssel előállított  $(\text{Ni}_{1-x}\text{Cu}_x)_{80}\text{P}_{20}$  amorf ötvözeteken ( $0 \leq x \leq 0,77$ ). A mért térfüggetlen második momentumok egyeznek a dipól-kiszélesedésre vonatkozó korábbi modellszámolásokkal. A P-P dipól kölcsönhatást a Cu tartalomtól függetlennek találtuk. Közlünk néhány adatot a Knight shiftre és a spin-rács relaxációs időre vonatkozólag is, rövid tárgyalással együtt.



## ABSTRACT

The  $^{31}\text{P}$  nuclear magnetic resonance was studied at room temperature as a function of the Ni/Cu ratio in rapidly quenched  $(\text{Ni}_{1-x}\text{Cu}_x)_80\text{P}_{20}$  ( $0 \leq x \leq 0.77$ ) amorphous alloys. The measured field-independent second moments agree with the results of earlier model calculations on dipolar broadening. The P-P dipolar interaction was found to be independent of the Cu content. Some data on Knight shift and spin-lattice relaxation time are also given and discussed briefly.

## INTRODUCTION

In this paper a detailed  $^{31}\text{P}$  nuclear magnetic resonance (NMR) study is reported on the rapidly quenched  $(\text{Ni}_{1-x}\text{Cu}_x)_80\text{P}_{20}$  amorphous alloy system for a wide range of the Ni/Cu ratio. The significance of the substitution of Ni by Cu atoms is twofold. The copper atoms have closed shell and, therefore, the replacement of Ni by Cu atoms is expected to cause considerable changes in the ratio of the s and d character of the conduction electrons at the Fermi surface in contrast to the Ni to Pd or Pt substitution [1] where such modifications are not expected. On the other hand, copper nuclei have high magnetic moment, thus they give rise to strong dipolar interactions with the  $^{31}\text{P}$  nuclei and this, in turn, results in a relatively high field-independent line-broadening effect, as reported already briefly [2], with respect to the Ni-P [2], Ni-Pd-P, and Ni-Pt-P [1] systems. In the present paper the attention will be focused on the field-independent line-broadening which is de-



terminated primarily by the geometrical arrangements of the atoms and the experimental second moments will be compared with the results of earlier model calculations [3]. Those of the  $^{31}\text{P}$  NMR parameters which reflect merely the electronic structure will be discussed only briefly here.

## SECOND MOMENT CALCULATIONS

The field-independent line-broadening of the  $^{31}\text{P}$  NMR spectrum originates from dipole-dipole interactions (since  $I = 1/2$  for  $^{31}\text{P}$  nuclei, there is no quadrupole interaction). The primary source is the direct dipolar coupling of the nuclear spins which can be further enhanced by indirect mechanisms (pseudo-exchange [4] or pseudo-dipolar [5] interactions) communicated via the conduction electrons. The magnitude of these pseudo-interactions is completely unknown in the present case. On the other hand, the direct dipolar interaction can be calculated by Van Vleck's theory [6] if the atomic coordinates are given. The total direct dipolar second moment  $M_2^D$  of the  $^{31}\text{P}$  NMR spectrum is given by

$$M_2^D = M_2(\text{P}) + \sum_i M_2(X_i) \quad (1)$$

where  $M_2(\text{P})$  is the second moment contribution due to the like  $^{31}\text{P}$  spins and  $M_2(X_i)$  is that due to the additional, unlike magnetic nuclei ( $X_i = {}^{61}\text{Ni}$ ,  ${}^{63}\text{Cu}$  and  ${}^{65}\text{Cu}$  in the present case). The second moment calculations were performed on a DRPHS model cluster generated by a Monte Carlo procedure. The following assumptions have been made: a) P-P nearest neighbours were not allowed; b) Ni and Cu are structurally equivalent having equal diameters and random substitution. It was found that  $M_2(\text{Ni}) \approx 0.02 M_2(\text{P})$  and, therefore, the contribution of the nickel nuclei can be neglected. Further details of the model calculations are described elsewhere [3].

## EXPERIMENTAL

The amorphous  $(\text{Ni}_{1-x}\text{Cu}_x)_{80}\text{P}_{20}$  alloy were prepared by rapid quenching from the melt using the Liebermann-Graham technique [7]. Glassy alloys were produced in the concentration range  $0 \leq x \leq 0.77$ .



The  $^{31}\text{P}$  NMR measurements were performed on a home-built continuous wave (CW) spectrometer and on a Bruker SXP 4-100 pulse spectrometer.

The field-dependence of the CW absorption derivative peak-to-peak linewidth  $\delta H$  could be well fitted by the formula [8]

$$(\delta H)^2 = (\delta H_0)^2 + (k_1 \cdot H)^2 \quad (2)$$

where  $\delta H_0$  is the field-independent linewidth contribution, the parameter  $k_1$  characterizes the strength of the field-dependence and  $H$  is the external magnetic field. As discussed in Ref. [3] the line shapes are very close to the Gaussian and therefore, the approximation

$$M_2^G = (\delta H_0)^2 / 4 \quad (3)$$

was used to obtain experimental field-independent second moments from the extrapolated  $\delta H_0$  values. The extrapolation procedure may result in large relative errors for  $M_2^G$  if  $k_1$  is high and  $\delta H_0$  is low. This is the case in Ni-P alloys but alloying with Cu changes the parameters favourably and makes the extrapolation reliable.

The second moment  $M_2^G$  includes both P-P and P-Cu interactions. On the other hand, with the help of coherent averaging techniques one can measure interactions between  $^{31}\text{P}$  nuclei separately. The P-P dipolar second moment  $M_2^{\text{P-P}}$  was measured by the two-pulse Carr-Purcell method described in details in Ref. [9].

## RESULTS AND DISCUSSION

The measured and calculated second moments of the  $^{31}\text{P}$  NMR spectrum are plotted on Fig. 1 as a function of the Cu content in amorphous  $(\text{Ni}_{1-x}\text{Cu}_x)_80\text{P}_{20}$  alloys. The calculated P-P second moment  $M_2(^{31}\text{P})$  is  $0.22 \text{ Oe}^2$ , and, by assumption, is independent of the Cu content. It can be seen that the second moment  $M_2^{\text{P-P}}$  is also independent of the Cu content but its  $0.4 \text{ Oe}^2$  value is almost twice the calculated one. This excess second moment is definitely greater than the uncertainty of the measured and calculated quantities. It can be concluded therefore, that  $M_2^{\text{P-P}}$  contains, besides the direct dipolar term  $M_2(^{31}\text{P})$ , contributions originating from



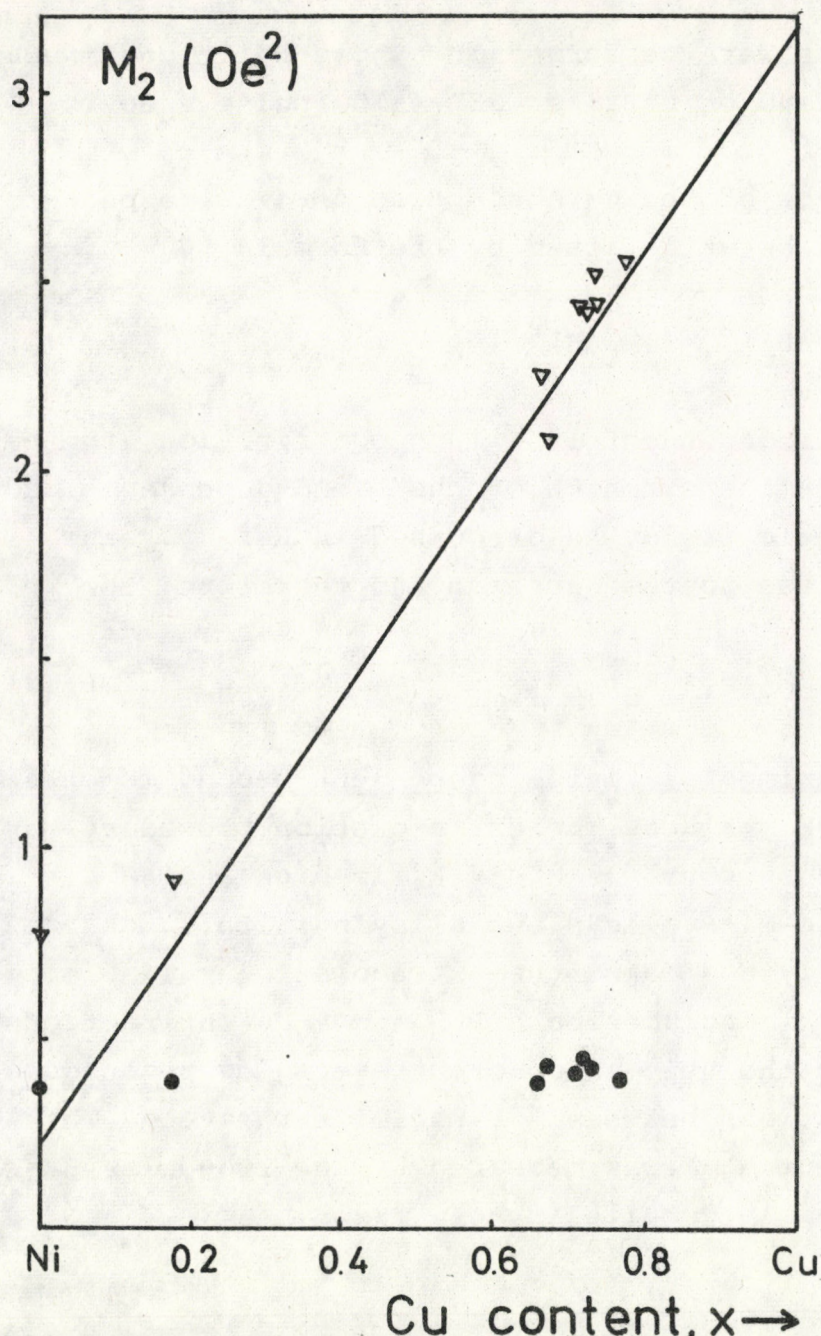


Fig. 1. Field-independent second moments of the  $^{31}\text{P}$  NMR spectrum vs Cu content in amorphous  $[\text{Ni}_{1-x}\text{Cu}_x]_{80}\text{P}_{20}$  alloys. Experimental:  $M_2^G/\nabla/$  and  $M_2^{P-P}/\bullet/$ ; calculated:  $M_2^D$  /straight line/ /see text for details/.

lated direct dipolar second moment  $M_2^D$  containing the contributions as given by eq.(1). In copper containing alloys the experimental data for  $M_2^G$  agree with the calculated values within their error, even if the observed P-P indirect interaction is taken into

indirect interactions, the magnitude of which seems to be independent of the Cu content and comparable to the direct dipolar term. On the other hand, other  $^{31}\text{P}$  NMR parameters such as Knight shift and spin-lattice relaxation time are sensitive to the amount of copper introduced into the Ni-P system as it will be given later and therefore, reflect considerable changes in the electronic structure upon alloying. However, in lack of more data on pseudo-interactions in amorphous alloys one cannot conclude at present regarding the origin of the excess second moment observed here.

The straight line on Fig. 1 represents the calcu-



account. Data point  $\nabla$  for  $x = 0$  may be left out of consideration since it has a high relative error for reasons given in the experimental section.

As discussed at some more length in Ref. [3], pseudo-interactions in amorphous Ni-Cu-P alloys, at least for high enough Cu content, are expected to be very low with respect to the direct dipolar broadening  $M_2^D$ . By considering *Fig. 1*, it means that DRPHS model clusters can be effectively used in estimating dipolar broadening in amorphous alloys. One can also conclude that Ni and Cu atoms are randomly distributed on the transition metal sites.

As to the low values of  $x$ , it is claimed by Durand et al. [10] that some metals, among them Cu too, do not substitute randomly for the matrix metal atoms in the dilute limit, but behave as glass former and substitute for it. This would mean that there will be no P-Cu first neighbours which, in turn, would considerably decrease the second moment contribution  $M_2(\text{Cu})$  and therefore, the total field-independent second moment of the  $^{31}\text{P}$  NMR spectrum. The only composition with low  $x$  on *Fig. 1* does not allow us to draw any conclusion in this respect, but more accurate measurements on amorphous  $(\text{Ni}_{1-x}\text{Cu}_x)_{80}\text{P}_{20}$  alloys with lower Cu content, with the help of some more sophisticated methods [9], will perhaps shed more light to this problem.

Finally, we briefly report on the room temperature measurements of the Knight shift  $K$  and the spin-lattice relaxation time  $T_1$ . It was found that upon introducing copper into the Ni-P system,  $K$  decreases and  $T_1$  increases with increasing Cu content. In a sample with  $x = 0.73$ , the spin-lattice relaxation time was measured between  $T = 100$  K and 300 K and the relation  $T_1 T = \text{const.}$  was obtained with  $T_1 T = 1.43$  K·s. The Korringa ratio  $k = K^2 T_1 T / S$ , where  $S = 1.605 \cdot 10^{-6}$  K·s for  $^{31}\text{P}$  nuclei, was found to change by a factor of two at room temperature in the concentration range studied and approached unity for high Cu content.

The significance of this result can be understood by taking into account that  $k = 1$  is obtained only if non-interacting s-type conduction electrons contribute to the Knight shift and the spin-lattice relaxation [11]. Detailed measurements of  $K$  and  $T_1$  will be published elsewhere.



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